the printed reproduction and is much smaller than the experimental uncertainties. The changes would be visible for the lower part of Fig. 1. The corrected full curve above that for YLAM is practically indistinguishable from the latter between 30° and 35°. Between 15° and 25° the full curve is too high by about 0.002 and at 30° by about 0.001(5). The inadvertent omission of reference to previous work on magnetic moment effects by Ohnuma[1] is acknowledged with apologies.

[1] Shoroku Ohnuma, Phys. Rev. **108**, 460 (1957). The calculations in this reference do not take into account wave distortion.

Low-Energy Nuclear Level Scheme of Rh^{104}, R. C. Greenwood [Phys. Rev. **129**, 345 (1963)]. The gamma-ray yield of the $B^{10}(n,\alpha)Li^7$ reaction was incorrectly given as 89.7% as well as being incorrectly referenced. The last sentence on p. 347 should therefore read: "In this boron spectrum, the 478-keV gamma rays are produced in 93.5% of all the resulting $B^{10}(n,\alpha)Li^7$ reactions,[9]" with Ref. 9 corrected to read: J. A. DeJuren and H. Rosenwasser, Phys. Rev. **93**, 831 (1954).

Influence of a Combined Magnetic Dipole and Electric Quadrupole Interaction on Angular Correlations, Kurt Alder, Eckart Matthias, Werner Schneider, and Rolf M. Steffen [Phys. Rev. **129**, 1199 (1963)]. Our expression for the anisotropy A [Eq. (57)] should be replaced by

$$A=[60A_{22}a_{22}^{(2)}-30\sqrt{3}(A_{24}+A_{42})a_{24}^{(2)}+45A_{44}a_{44}^{(2)}]\times\Big[16(2I+1)+10A_{22}(a_{22}^{(0)}-3a_{22}^{(2)})+\tfrac{9}{2}A_{44}\Big(\frac{9}{4}a_{44}^{(0)}-5a_{44}^{(2)}+\frac{35}{4}a_{44}^{(4)}\Big)-\tfrac{3}{2}(A_{24}+A_{42})(3(5)^{1/2}a_{24}^{0}-10\sqrt{3}a_{24}^{(2)})\Big]^{-1}.$$

Therefore, Figs. 18 and 19 have to be disregarded and replaced by the following figures:

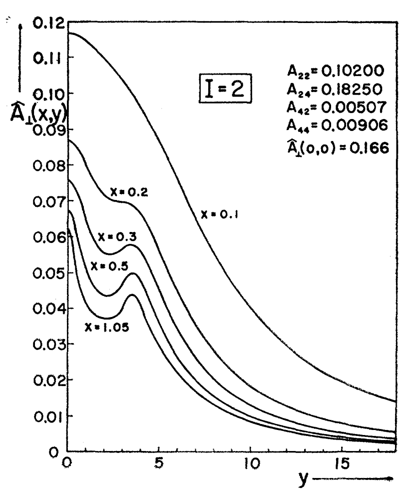

Fig. 18. The anisotropy $\hat{A}_\perp=[W(\pi)-W(\frac{1}{2}\pi)]/W(\frac{1}{2}\pi)$ of the integral angular correlation with magnetic field perpendicular to detector plane for $I=2$.

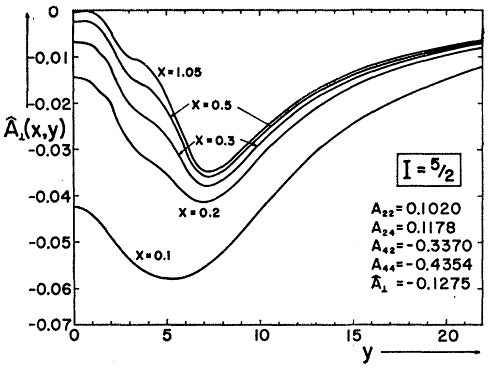

Fig. 19. The anisotropy $\hat{A}_\perp=[W(\pi)-W(\frac{1}{2}\pi)]/W(\frac{1}{2}\pi)$ of the integral angular correlation with magnetic field perpendicular to detector plane for $I=\frac{5}{2}$.

In addition, Eq. (61) should read

$$W(\theta)=\sum_{k_1k_2}A_{k_1}(R_1)A_{k_2}(R_2)\times[(2k_1+1)(2k_2+1)]^{1/2}a_{k_1k_2}{}^{0}P_k(\cos\theta),$$

where $k=k_1$ if $\theta_1=\theta$, and $k=k_2$ if $\theta_2=\theta$. In the following two sentences $a_{kk}{}^0$ and $a_{kk}{}^0$ should be replaced by $a_{k_1k_2}{}^0$ and $a_{k_1k_2}{}^0$, respectively.

Angular Correlation Perturbed by an Anisotropic Hyperfine Interaction. H. J. Leisi and R. T. Deck [Phys. Rev. **129**, 2117 (1963)]. In the final version of the manuscript a factor $1/(4\pi)^{1/2}$ was omitted from Eqs. (31) and (32). These equations should read

$$Y_k{}^\mu(\theta,\phi)=\frac{a_k{}^\mu(\theta)}{(4\pi)^{1/2}}e^{i\mu\phi},\tag{31}$$

$$Y_k{}^\mu(0,\phi)=\frac{a_k{}^\mu(0)}{(4\pi)^{1/2}}e^{i\mu\phi}=\Big(\frac{2k+1}{4\pi}\Big)^{1/2}\delta_{\mu0}.\tag{32}$$

No figures or other formulas are affected by the correction. In the last sentence of Sec. VI the phrase "parallel to the detector plane" should read "perpendicular to the detector plane."

Properties and Effects of η Decays. Riazuddin and Fayyazuddin [Phys. Rev. **129**, 2337 (1963)]. Due to use of a normalization different from that of Chew, the value of $\lambda/4\pi=-0.15$ used in the text should be replaced by $\lambda/16\pi=-0.15$. Then Eqs. (9) and (10) of the text, respectively, become

$$\Gamma_\eta(\pi^+\pi^-\pi^0)\approx224\text{ eV},\tag{9}$$

$$\Gamma_\eta(3\pi^0)\approx358\text{ eV}.\tag{10}$$

The conclusions after Eq. (11) in the first and second paragraphs should read as follows: "Combining the estimate of Hori *et al.* for $\Gamma_\eta(2\gamma)$ or the estimate $\Gamma_\eta(2\gamma)\approx192$ eV with our estimates (9) and (10) for $\Gamma_\eta(\pi^+\pi^-\pi^0)$ and $\Gamma_\eta(3\pi^0)$, we find $\Gamma_\eta(2\gamma)$ different from $\Gamma_\eta(3\pi^0)$ and $r=\Gamma_\eta(\text{neutrals})/\Gamma_\eta(\pi^+\pi^-\pi^0)\approx2.4$

if we use $\Gamma_\eta(2\gamma)=192$ eV, and $r\approx 1.7$ when we employ $\Gamma_\eta(2\gamma)=25$ eV. The ratio 2.4 is consistent with experiment. Estimates for $\Gamma_\eta(2\gamma)$ made on a unitary symmetry model also lie between 25 eV and 192 eV and, hence, for these estimates r will lie between 1.7 and 2.4. Hori *et al.* as well as Gell-Mann *et al.* $\cdots$."

If $(\delta g_{\pi NN}/g_{\pi NN})$ is taken to be 0.7%, then the values (9) and (10) for $\Gamma_\eta(\pi^+\pi^-\pi^0)$ and $\Gamma_\eta(3\pi^0)$ are unchanged provided that $g_{\eta NN}{}^2/4\pi\approx 1$. However, then Eq. (11) gives $\Gamma_\eta(2\gamma)\approx 12$ eV so that r lies between 1.6 and 2.4 according as $\Gamma_\eta(2\gamma)$ lies between 12 eV and 192 eV. Again with $(\delta g_{\pi NN}/g_{\pi NN})\approx 0.7\%$ and $g_{\eta NN}{}^2/4\pi\approx 2$, $\Gamma_\eta(\pi^+\pi^-\pi^0)$ and $\Gamma_\eta(3\pi^0)$ become, respectively, 112 eV and 179 eV. These values are consistent with the estimates of Barret and Barton[1] who estimated the $\eta^0\to\pi^0$ vertex by a quite different approach based on unitary symmetry. In this case r lies between 1.8 and 3.3 according as $\Gamma_\eta(2\gamma)$ lies between 25 eV and 192 eV, consistent with experiment. In this case $\Gamma_\eta(2\gamma)$ can be equal to $\Gamma_\eta(3\pi^0)$ depending on what value we take between 25 eV and 192 eV for $\Gamma_\eta(2\gamma)$.

In the last paragraph but one, if we take $\lambda/16\pi=-0.15$, the pseudoscalar coupling constant $g_{\Sigma NK}{}^2/4\pi$ has the values 48 to 24 according as $R(K_2{}^0\to\pi^+\pi^-\pi^0)=1.5\times 10^6$ sec^{-1} or 3×10^6 sec^{-1}. For scalar $K\Sigma N$ coupling,

$$g_{\Sigma NK}{}^2/4\pi\approx 0.48.$$

Thus, for pseudoscalar coupling, $g_{\Sigma NK}{}^2/4\pi$ comes out to be quite large, showing that the K pole in $\Sigma^-\to n+\pi^-$ does not dominate and that one has to consider other contributions also.

We are grateful to Barbara Barrett for pointing out the error in our paper.

[1] Barbara Barret and G. Barton (to be published).

Angular Distribution of Muons in π-μ Decay at Rest, H. Hulubei, J. S. Ausländer, E. M. Friedländer, and Ş. Ţiţeica [Phys. Rev. **129**, 2789 (1963)]. 1. In Table I, column headed "Sample size," row "Ω^*": instead of 19126^b read (19126)^b. This figure does not represent an *actual* sample size, but a *fictitious* one. 2. In Fig. 8, (a) and (b) must be interchanged in order to obtain agreement between (i) drawings and (ii) figure caption and text.

Relaxation-Time Measurements in Ruby by a dc Magnetization Technique, Shih-Yu Feng and N. Bloembergen [Phys. Rev. **130**, 531 (1963)]. In the caption of Fig. 4 and in the line of the text immediately following this figure, it is erroneously stated that "$H_{dc}=2990$ G." This should read "$H_{dc}=1580\pm 20$ G." The value originally quoted belongs to another transition at 0^0 orientation. A check of our experimental records revealed the correct value, although the precision is rather poor. An accurate machine solution of the spin Hamiltonian at the frequency used in the experiment gives the following result for the harmonic point:

Orientation	Resonance H_{dc} (G)	Ratio ν_{34}/ν_{13}
21°	1585	2.92:2
22°	1674	3.01:2

We wish to thank Dr. W. Grant for calling our attention to this error.

Partial-Wave Bethe-Salpeter Equation, Noboru Nakanishi [Phys. Rev. **130**, 1230 (1963)]. In the denominators of Eqs. (3.21) and (4.5), and in the argument of the δ function of Eq. (4.9), x and $(1-x)$ should be interchanged.

Branching Ratios of π^- Mesons Stopped in Hydrogen and Deuterium, James W. Ryan [Phys. Rev. **130**, 1554 (1963)]. Add:

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